

# Entropy and mixing for $\mathbb{Z}^d$ SFTs

Ronnie Pavlov

University of Denver  
[www.math.du.edu/~rpavlov](http://www.math.du.edu/~rpavlov)

1st School on Dynamical Systems and Computation (DySyCo)  
CMM, Santiago, Chile

# Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$

## Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$
- Any such  $\mu$  is determined by values on cylinder sets  $[w]$

## Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$
- Any such  $\mu$  is determined by values on cylinder sets  $[w]$
- To any such measure is assigned measure-theoretic entropy:

## Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$
- Any such  $\mu$  is determined by values on cylinder sets  $[w]$
- To any such measure is assigned measure-theoretic entropy:

$$h(\mu) = \lim_{n \rightarrow \infty} \frac{-1}{n^d} \sum_{w \in L(X) \cap A^{\{1, \dots, n\}^d}} \mu(w) \log \mu(w)$$

## Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$
- Any such  $\mu$  is determined by values on cylinder sets  $[w]$
- To any such measure is assigned measure-theoretic entropy:

$$h(\mu) = \lim_{n \rightarrow \infty} \frac{-1}{n^d} \sum_{w \in L(X) \cap A^{\{1, \dots, n\}^d}} \mu(w) \log \mu(w)$$

- Note: if  $\mu$  uniformly distributed over patterns in  $L(X) \cap A^{\{1, \dots, n\}^d}$ :

# Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$
- Any such  $\mu$  is determined by values on cylinder sets  $[w]$
- To any such measure is assigned measure-theoretic entropy:

$$h(\mu) = \lim_{n \rightarrow \infty} \frac{-1}{n^d} \sum_{w \in L(X) \cap A^{\{1, \dots, n\}^d}} \mu(w) \log \mu(w)$$

- Note: if  $\mu$  uniformly distributed over patterns in  $L(X) \cap A^{\{1, \dots, n\}^d}$ :

- $h(\mu) = \lim_{n \rightarrow \infty} \frac{-1}{n^d} \log \left( \frac{1}{|L(X) \cap A^{\{1, \dots, n\}^d}|} \right)$

## Measures on $\mathbb{Z}^d$ subshifts

- All measures we consider will be shift-invariant probability Borel measures on  $A^{\mathbb{Z}^d}$
- Any such  $\mu$  is determined by values on cylinder sets  $[w]$
- To any such measure is assigned measure-theoretic entropy:

$$h(\mu) = \lim_{n \rightarrow \infty} \frac{-1}{n^d} \sum_{w \in L(X) \cap A^{\{1, \dots, n\}^d}} \mu(w) \log \mu(w)$$

- Note: if  $\mu$  uniformly distributed over patterns in  $L(X) \cap A^{\{1, \dots, n\}^d}$ :

$$\bullet h(\mu) = \lim_{n \rightarrow \infty} \frac{-1}{n^d} \log \left( \frac{1}{|L(X) \cap A^{\{1, \dots, n\}^d}|} \right) = h(X)$$



# Measures on $\mathbb{Z}^d$ subshifts

- There usually does not exist such a uniformly distributed measure (PROVE)

# Measures on $\mathbb{Z}^d$ subshifts

- There usually does not exist such a uniformly distributed measure (PROVE)
- Nevertheless, we have the classical Variational Principle:

# Measures on $\mathbb{Z}^d$ subshifts

- There usually does not exist such a uniformly distributed measure (PROVE)
- Nevertheless, we have the classical Variational Principle:
- **Theorem: (Variational Principle)**  $\sup h(\mu) = h(X)$  (over  $\mu$  with support in  $X$ ), and the sup is achieved

## Measures on $\mathbb{Z}^d$ subshifts

- There usually does not exist such a uniformly distributed measure (PROVE)
- Nevertheless, we have the classical Variational Principle:
- **Theorem: (Variational Principle)**  $\sup h(\mu) = h(X)$  (over  $\mu$  with support in  $X$ ), and the sup is achieved
- Measures  $\mu$  for which  $h(\mu) = h(X)$  are called **measures of maximal entropy**

## Measures on $\mathbb{Z}^d$ subshifts

- There usually does not exist such a uniformly distributed measure (PROVE)
- Nevertheless, we have the classical Variational Principle:
- **Theorem: (Variational Principle)**  $\sup h(\mu) = h(X)$  (over  $\mu$  with support in  $X$ ), and the sup is achieved
- Measures  $\mu$  for which  $h(\mu) = h(X)$  are called **measures of maximal entropy**
- Such measures are useful for studying topological entropy, since they allow the additional strength of measure theory to be brought to bear

# Measures of maximal entropy

- Any measure of maximal entropy  $\mu$  for an SFT  $X$  has an interesting property

# Measures of maximal entropy

- Any measure of maximal entropy  $\mu$  for an SFT  $X$  has an interesting property
- **Theorem:** (Burton-Steif/Lanford-Ruelle) For any such  $\mu$ , any finite  $S$  and  $T \supset \partial S$  for which  $S \cap T = \emptyset$ , and for any  $\delta \in L_T(X)$ ,  $\mu(x|_S : x|_T = \delta)$  is uniform over all  $x \in L_S(X)$  for which  $x\delta \in L(X)$ .

# Measures of maximal entropy

- Any measure of maximal entropy  $\mu$  for an SFT  $X$  has an interesting property
- **Theorem:** (Burton-Steif/Lanford-Ruelle) For any such  $\mu$ , any finite  $S$  and  $T \supset \partial S$  for which  $S \cap T = \emptyset$ , and for any  $\delta \in L_T(X)$ ,  $\mu(x|_S : x|_T = \delta)$  is uniform over all  $x \in L_S(X)$  for which  $x\delta \in L(X)$ .
  - Call such measures **uniform Gibbs measures**.



# Measures of maximal entropy

- Any measure of maximal entropy  $\mu$  for an SFT  $X$  has an interesting property
- **Theorem:** (Burton-Steif/Lanford-Ruelle) For any such  $\mu$ , any finite  $S$  and  $T \supset \partial S$  for which  $S \cap T = \emptyset$ , and for any  $\delta \in L_T(X)$ ,  $\mu(x|_S : x|_T = \delta)$  is uniform over all  $x \in L_S(X)$  for which  $x\delta \in L(X)$ .
  - Call such measures **uniform Gibbs measures**.
- Example:  $\mathcal{H}$  the  $\mathbb{Z}^2$  hard square shift: if  $\mu$  is uniform Gibbs,

# Measures of maximal entropy

- Any measure of maximal entropy  $\mu$  for an SFT  $X$  has an interesting property
- **Theorem:** (Burton-Steif/Lanford-Ruelle) For any such  $\mu$ , any finite  $S$  and  $T \supset \partial S$  for which  $S \cap T = \emptyset$ , and for any  $\delta \in L_T(X)$ ,  $\mu(x|_S : x|_T = \delta)$  is uniform over all  $x \in L_S(X)$  for which  $x\delta \in L(X)$ .
  - Call such measures **uniform Gibbs measures**.
- Example:  $\mathcal{H}$  the  $\mathbb{Z}^2$  hard square shift: if  $\mu$  is uniform Gibbs, conditioned on  $\begin{matrix} 1 & 0 & 1 & 0 \\ 0 & & 0 & \\ 0 & & 1 & \\ 1 & 0 & 0 & 0 \end{matrix}$ , fillings  $\begin{matrix} 0 & 0 \\ 0 & 0 \end{matrix}$ ,  $\begin{matrix} 0 & 0 \\ 1 & 0 \end{matrix}$ ,  $\begin{matrix} 1 & 0 \\ 0 & 0 \end{matrix}$  equally probable.

# Measures of maximal entropy

- Any measure of maximal entropy  $\mu$  for an SFT  $X$  has an interesting property
- Theorem:** (Burton-Steif/Lanford-Ruelle) For any such  $\mu$ , any finite  $S$  and  $T \supset \partial S$  for which  $S \cap T = \emptyset$ , and for any  $\delta \in L_T(X)$ ,  $\mu(x|_S : x|_T = \delta)$  is uniform over all  $x \in L_S(X)$  for which  $x\delta \in L(X)$ .

- Call such measures **uniform Gibbs measures**.

- Example:  $\mathcal{H}$  the  $\mathbb{Z}^2$  hard square shift: if  $\mu$  is uniform Gibbs,

conditioned on  $\begin{matrix} 1 & 0 & 1 & 0 \\ 0 & & 0 & \\ 0 & & 1 & \\ 1 & 0 & 0 & 0 \end{matrix}$ , fillings  $\begin{matrix} 0 & 0 \\ 0 & 0 \end{matrix}$ ,  $\begin{matrix} 0 & 0 \\ 1 & 0 \end{matrix}$ ,  $\begin{matrix} 1 & 0 \\ 0 & 0 \end{matrix}$  equally probable.

- Same conditional probabilities if  $\begin{matrix} 1 & 0 & 1 & 0 \\ 0 & & 0 & \\ 0 & & 1 & \\ 1 & 0 & 0 & 0 \end{matrix}$  changed to  $\begin{matrix} & 0 & 1 & 0 & 0 & 1 \\ 0 & 1 & 0 & 1 & 0 & \\ 1 & 0 & & 0 & 0 & \\ 0 & 0 & & 1 & & \\ 0 & 1 & 0 & 0 & 0 & 0 \\ 1 & 0 & & & & 1 \end{matrix}$

# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT

# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT
- One way to create a uniform Gibbs measure is as weak limit of uniform measures on finite sets given boundary conditions (PROVE)

# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT
- One way to create a uniform Gibbs measure is as weak limit of uniform measures on finite sets given boundary conditions (PROVE)
- An SFT can have multiple uniform Gibbs measures

# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT
- One way to create a uniform Gibbs measure is as weak limit of uniform measures on finite sets given boundary conditions (PROVE)
- An SFT can have multiple uniform Gibbs measures
- Easy example:  $X = \{0\}^{\mathbb{Z}^2} \cup \{1\}^{\mathbb{Z}^2}$ .

# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT
- One way to create a uniform Gibbs measure is as weak limit of uniform measures on finite sets given boundary conditions (PROVE)
- An SFT can have multiple uniform Gibbs measures
- Easy example:  $X = \{0\}^{\mathbb{Z}^2} \cup \{1\}^{\mathbb{Z}^2}$ .
  - Both the measure supported on the point  $\{0\}^{\mathbb{Z}^d}$  and the measure supported on  $\{1\}^{\mathbb{Z}^2}$  are uniform Gibbs measures



# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT
- One way to create a uniform Gibbs measure is as weak limit of uniform measures on finite sets given boundary conditions (PROVE)
- An SFT can have multiple uniform Gibbs measures
- Easy example:  $X = \{0\}^{\mathbb{Z}^2} \cup \{1\}^{\mathbb{Z}^2}$ .
  - Both the measure supported on the point  $\{0\}^{\mathbb{Z}^d}$  and the measure supported on  $\{1\}^{\mathbb{Z}^2}$  are uniform Gibbs measures
  - Reason for non-uniqueness: boundary conditions of all 0 and all 1 on large finite sets induce drastically different uniform measures on interior

# Measures of maximal entropy

- Uniform Gibbs measures are “as uniform as possible” measures on an SFT
- One way to create a uniform Gibbs measure is as weak limit of uniform measures on finite sets given boundary conditions (PROVE)
- An SFT can have multiple uniform Gibbs measures
- Easy example:  $X = \{0\}^{\mathbb{Z}^2} \cup \{1\}^{\mathbb{Z}^2}$ .
  - Both the measure supported on the point  $\{0\}^{\mathbb{Z}^d}$  and the measure supported on  $\{1\}^{\mathbb{Z}^2}$  are uniform Gibbs measures
  - Reason for non-uniqueness: boundary conditions of all 0 and all 1 on large finite sets induce drastically different uniform measures on interior
  - We'll return to this

## Consequences of topological mixing conditions

- **Theorem:** (Dobrushin) If  $X$  is a strongly irreducible SFT, then any (shift-invariant) uniform Gibbs measure is also a measure of maximal entropy.

# Consequences of topological mixing conditions

- **Theorem:** (Dobrushin) If  $X$  is a strongly irreducible SFT, then any (shift-invariant) uniform Gibbs measure is also a measure of maximal entropy.
  - In fact Dobrushin used an even weaker condition, which we don't state here

## Consequences of topological mixing conditions

- **Theorem:** (Dobrushin) If  $X$  is a strongly irreducible SFT, then any (shift-invariant) uniform Gibbs measure is also a measure of maximal entropy.
  - In fact Dobrushin used an even weaker condition, which we don't state here
- In addition, measures of maximal entropy on strongly irreducible SFTs must be fully supported (PROVE)

## Consequences of topological mixing conditions

- **Theorem:** (Dobrushin) If  $X$  is a strongly irreducible SFT, then any (shift-invariant) uniform Gibbs measure is also a measure of maximal entropy.
  - In fact Dobrushin used an even weaker condition, which we don't state here
- In addition, measures of maximal entropy on strongly irreducible SFTs must be fully supported (PROVE)
- This implies the claim from Lecture 1 that strongly irreducible SFTs are entropy minimal (have no proper subshifts with equal entropy)

## Consequences of topological mixing conditions

- If  $X$  is block gluing, uniform Gibbs measures are not necessarily measures of maximal entropy

# Consequences of topological mixing conditions

- If  $X$  is block gluing, uniform Gibbs measures are not necessarily measures of maximal entropy
- Consider the southeast shift defined in Lecture 1 (no  $\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$ )



## Consequences of topological mixing conditions

- If  $X$  is block gluing, uniform Gibbs measures are not necessarily measures of maximal entropy
- Consider the southeast shift defined in Lecture 1 (no  $\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$ )
- For this shift, the measure  $\mu$  supported on the single fixed point  $0^{\mathbb{Z}^2}$  is uniform Gibbs (comes from weak limit of boundary conditions of all 0s)

## Consequences of topological mixing conditions

- If  $X$  is block gluing, uniform Gibbs measures are not necessarily measures of maximal entropy
- Consider the southeast shift defined in Lecture 1 (no  $\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$ )
- For this shift, the measure  $\mu$  supported on the single fixed point  $0^{\mathbb{Z}^2}$  is uniform Gibbs (comes from weak limit of boundary conditions of all 0s)
- But then  $h(\mu) = 0$ , and  $h(S) > 0$ , so  $\mu$  not an m.m.e.

## Consequences of topological mixing conditions

- If  $X$  is block gluing, uniform Gibbs measures are not necessarily measures of maximal entropy
- Consider the southeast shift defined in Lecture 1 (no  $\begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$ )
- For this shift, the measure  $\mu$  supported on the single fixed point  $0^{\mathbb{Z}^2}$  is uniform Gibbs (comes from weak limit of boundary conditions of all 0s)
- But then  $h(\mu) = 0$ , and  $h(S) > 0$ , so  $\mu$  not an m.m.e.
- In addition, measures of maximal entropy on block gluing SFTs are not necessarily fully supported

# Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing

## Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing
- **iceberg model**  $\mathcal{I}_M$  (Burton-Steif):  $A = \{-M, \dots, -1, 1, \dots, M\}$ ,  $\mathcal{F} = \{\text{adjacent pairs } i, j \text{ with } ij < -1\}$ .

## Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing
- **iceberg model**  $\mathcal{I}_M$  (Burton-Steif):  $A = \{-M, \dots, -1, 1, \dots, M\}$ ,  $\mathcal{F} = \{\text{adjacent pairs } i, j \text{ with } ij < -1\}$ .
- Only allowed adjacent integers with opposite signs are  $\pm 1$ .

# Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing
- **iceberg model**  $\mathcal{I}_M$  (Burton-Steif):  $A = \{-M, \dots, -1, 1, \dots, M\}$ ,  $\mathcal{F} = \{\text{adjacent pairs } i, j \text{ with } ij < -1\}$ .
- Only allowed adjacent integers with opposite signs are  $\pm 1$ .
- $\mathcal{I}_M$  is strongly irreducible: can use  $\pm 1$  to mix between any two patterns (PROVE)

## Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing
- **iceberg model**  $\mathcal{I}_M$  (Burton-Steif):  $A = \{-M, \dots, -1, 1, \dots, M\}$ ,  $\mathcal{F} = \{\text{adjacent pairs } i, j \text{ with } ij < -1\}$ .
- Only allowed adjacent integers with opposite signs are  $\pm 1$ .
- $\mathcal{I}_M$  is strongly irreducible: can use  $\pm 1$  to mix between any two patterns (PROVE)
  - $\mu^+, \mu^-$  uniform Gibbs measures/MMEs obtained by weak limits of conditioning on boundaries of all  $M$ s, all  $-M$ s respectively



## Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing
- **iceberg model**  $\mathcal{I}_M$  (Burton-Steif):  $A = \{-M, \dots, -1, 1, \dots, M\}$ ,  $\mathcal{F} = \{\text{adjacent pairs } i, j \text{ with } ij < -1\}$ .
- Only allowed adjacent integers with opposite signs are  $\pm 1$ .
- $\mathcal{I}_M$  is strongly irreducible: can use  $\pm 1$  to mix between any two patterns (PROVE)
  - $\mu^+, \mu^-$  uniform Gibbs measures/MMEs obtained by weak limits of conditioning on boundaries of all  $M$ s, all  $-M$ s respectively
  - For large  $M$ ,  $\mu^+ \neq \mu^-$  (will prove in Lecture 4)

## Multiple measures of maximal entropy

- There are examples with extremely strong topological mixing properties whose uniform Gibbs measures are quite nonmixing
- **iceberg model**  $\mathcal{I}_M$  (Burton-Steif):  $A = \{-M, \dots, -1, 1, \dots, M\}$ ,  $\mathcal{F} = \{\text{adjacent pairs } i, j \text{ with } ij < -1\}$ .
- Only allowed adjacent integers with opposite signs are  $\pm 1$ .
- $\mathcal{I}_M$  is strongly irreducible: can use  $\pm 1$  to mix between any two patterns (PROVE)
  - $\mu^+, \mu^-$  uniform Gibbs measures/MMEs obtained by weak limits of conditioning on boundaries of all  $M$ s, all  $-M$ s respectively
  - For large  $M$ ,  $\mu^+ \neq \mu^-$  (will prove in Lecture 4)
  - It's possible to transition from a positive boundary condition to a negative letter at 0 (strong irreducibility), but very unlikely (multiple uniform Gibbs measures)

# Mixing for measures of maximal entropy

- When there is a unique uniform Gibbs measure  $\mu$ , this means that “boundary influence decays with distance”

# Mixing for measures of maximal entropy

- When there is a unique uniform Gibbs measure  $\mu$ , this means that “boundary influence decays with distance”
- This is a type of measure-theoretic mixing, called **weak spatial mixing**

# Mixing for measures of maximal entropy

- When there is a unique uniform Gibbs measure  $\mu$ , this means that “boundary influence decays with distance”
- This is a type of measure-theoretic mixing, called **weak spatial mixing**
- Formally, for a function  $f(n) \rightarrow 0$ , we say that  $\mu$  is **weak spatial mixing with rate**  $f(n)$  if for every  $n \in \mathbb{N}$ ,  $a \in A$ , finite set  $T \supseteq \{-n, \dots, n\}^d$ , and  $\delta, \delta' \in A^{\partial T}$ ,

# Mixing for measures of maximal entropy

- When there is a unique uniform Gibbs measure  $\mu$ , this means that “boundary influence decays with distance”
- This is a type of measure-theoretic mixing, called **weak spatial mixing**
- Formally, for a function  $f(n) \rightarrow 0$ , we say that  $\mu$  is **weak spatial mixing with rate  $f(n)$**  if for every  $n \in \mathbb{N}$ ,  $a \in A$ , finite set  $T \supseteq \{-n, \dots, n\}^d$ , and  $\delta, \delta' \in A^{\partial T}$ ,

$$|\mu(x(0) = a \mid x(\partial T) = \delta) - \mu(x(0) = a \mid x(\partial T) = \delta')| < f(n).$$

# Mixing for measures of maximal entropy

- When there is a unique uniform Gibbs measure  $\mu$ , this means that “boundary influence decays with distance”
- This is a type of measure-theoretic mixing, called **weak spatial mixing**
- Formally, for a function  $f(n) \rightarrow 0$ , we say that  $\mu$  is **weak spatial mixing with rate  $f(n)$**  if for every  $n \in \mathbb{N}$ ,  $a \in A$ , finite set  $T \supseteq \{-n, \dots, n\}^d$ , and  $\delta, \delta' \in A^{\partial T}$ ,  
$$|\mu(x(0) = a \mid x(\partial T) = \delta) - \mu(x(0) = a \mid x(\partial T) = \delta')| < f(n).$$
- Differs from usual notion of measure-theoretic mixing due to (necessary) idea of “surrounding”

# Mixing for measures of maximal entropy

- We say that  $\mu$  is **strongly spatial mixing with rate**  $f(n)$  if for any every  $n \in \mathbb{N}$ ,  $a \in A$ , finite set  $T \ni 0$ , and for any patterns  $\delta, \delta' \in A^{\partial T}$  agreeing for all  $v$  with  $\|v\| < n$ ,



# Mixing for measures of maximal entropy

- We say that  $\mu$  is **strongly spatial mixing with rate**  $f(n)$  if for any every  $n \in \mathbb{N}$ ,  $a \in A$ , finite set  $T \ni 0$ , and for any patterns  $\delta, \delta' \in A^{\partial T}$  agreeing for all  $v$  with  $\|v\| < n$ ,

$$|\mu(x(0) = a \mid x(\partial T) = \delta) - \mu(x(0) = a \mid x(\partial T) = \delta')| < f(n).$$

# Mixing for measures of maximal entropy

- We say that  $\mu$  is **strongly spatial mixing with rate**  $f(n)$  if for any every  $n \in \mathbb{N}$ ,  $a \in A$ , finite set  $T \ni 0$ , and for any patterns  $\delta, \delta' \in A^{\partial T}$  agreeing for all  $v$  with  $\|v\| < n$ ,

$$|\mu(x(0) = a \mid x(\partial T) = \delta) - \mu(x(0) = a \mid x(\partial T) = \delta')| < f(n).$$

- Different from WSM since we allow the boundary conditions  $\delta$  and  $\delta'$  to contain sites close to 0, as long as they agree on these sites.

# Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] \mid [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] \mid [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$
  - $\min_{\delta} \mu([w] \mid [\delta]) \leq \mu([w]) \leq \max_{\delta} \mu([w] \mid [\delta])$

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] \mid [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$
  - $\min_{\delta} \mu([w] \mid [\delta]) \leq \mu([w]) \leq \max_{\delta} \mu([w] \mid [\delta])$
  - Upper and lower bounds tighten with rate of WSM

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] \mid [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$
  - $\min_{\delta} \mu([w] \mid [\delta]) \leq \mu([w]) \leq \max_{\delta} \mu([w] \mid [\delta])$
  - Upper and lower bounds tighten with rate of WSM
- With SSM, can efficiently approximate conditional measures as well:



## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] | [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$
  - $\min_{\delta} \mu([w] | [\delta]) \leq \mu([w]) \leq \max_{\delta} \mu([w] | [\delta])$
  - Upper and lower bounds tighten with rate of WSM
- With SSM, can efficiently approximate conditional measures as well:
  - $\mu([w] | [v]) = \sum_{\delta} \mu([\delta] | [v])\mu([w] | [\delta], [v])$

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] \mid [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$
  - $\min_{\delta} \mu([w] \mid [\delta]) \leq \mu([w]) \leq \max_{\delta} \mu([w] \mid [\delta])$
  - Upper and lower bounds tighten with rate of WSM
- With SSM, can efficiently approximate conditional measures as well:
  - $\mu([w] \mid [v]) = \sum_{\delta} \mu([\delta] \mid [v])\mu([w] \mid [\delta], [v])$
  - $\min_{\delta} \mu([w] \mid [\delta], [v]) \leq \mu([w]) \leq \max_{\delta} \mu([w] \mid [\delta], [v])$

## Consequences of spatial mixing

- Spatial mixing properties are very useful and powerful
- With WSM, can at least approximate measures of cylinder sets:
  - For a fixed  $n$ ,  $\mu([w]) = \sum_{\delta} \mu([\delta])\mu([w] | [\delta])$ , where  $\delta \in A^{\partial\{-n, \dots, n\}^d}$
  - $\min_{\delta} \mu([w] | [\delta]) \leq \mu([w]) \leq \max_{\delta} \mu([w] | [\delta])$
  - Upper and lower bounds tighten with rate of WSM
- With SSM, can efficiently approximate conditional measures as well:
  - $\mu([w] | [v]) = \sum_{\delta} \mu([\delta] | [v])\mu([w] | [\delta], [v])$
  - $\min_{\delta} \mu([w] | [\delta], [v]) \leq \mu([w]) \leq \max_{\delta} \mu([w] | [\delta], [v])$
  - Upper and lower bounds tighten with rate of SSM

## Consequences of spatial mixing

- Such efficient approximations imply stronger computability properties of entropy

## Consequences of spatial mixing

- Such efficient approximations imply stronger computability properties of entropy
- **Theorem:** (Marcus-Pavlov) If  $X$  is a  $\mathbb{Z}^2$  SFT whose measure of maximal entropy has SSM with exponential rate, then  $h(X)$  is computable in polynomial time, i.e. you can get upper and lower approximations to within tolerance  $\frac{1}{n}$  in  $n^{O(1)}$  steps

## Consequences of spatial mixing

- Such efficient approximations imply stronger computability properties of entropy
- **Theorem:** (Marcus-Pavlov) If  $X$  is a  $\mathbb{Z}^2$  SFT whose measure of maximal entropy has SSM with exponential rate, then  $h(X)$  is computable in polynomial time, i.e. you can get upper and lower approximations to within tolerance  $\frac{1}{n}$  in  $n^{O(1)}$  steps
- Cliffhanger: The  $\mathbb{Z}^2$  hard square shift,  $\mathcal{H}$ , has a unique MME with SSM with exponential rate!

## Consequences of spatial mixing

- Such efficient approximations imply stronger computability properties of entropy
- **Theorem:** (Marcus-Pavlov) If  $X$  is a  $\mathbb{Z}^2$  SFT whose measure of maximal entropy has SSM with exponential rate, then  $h(X)$  is computable in polynomial time, i.e. you can get upper and lower approximations to within tolerance  $\frac{1}{n}$  in  $n^{O(1)}$  steps
- Cliffhanger: The  $\mathbb{Z}^2$  hard square shift,  $\mathcal{H}$ , has a unique MME with SSM with exponential rate!
- Corollary:  $h(\mathcal{H})$  is computable in polynomial time

## Consequences of spatial mixing

- Such efficient approximations imply stronger computability properties of entropy
- **Theorem:** (Marcus-Pavlov) If  $X$  is a  $\mathbb{Z}^2$  SFT whose measure of maximal entropy has SSM with exponential rate, then  $h(X)$  is computable in polynomial time, i.e. you can get upper and lower approximations to within tolerance  $\frac{1}{n}$  in  $n^{O(1)}$  steps
- Cliffhanger: The  $\mathbb{Z}^2$  hard square shift,  $\mathcal{H}$ , has a unique MME with SSM with exponential rate!
- Corollary:  $h(\mathcal{H})$  is computable in polynomial time
- More about this next time